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► To cite this version:

Denis Bonheure, Pietro D'Avenia, Alessio Pomponio. On the electrostatic Born-Infeld equation with extended charges. Communications in Mathematical Physics, 2016, 10.1007/s00220-016-2586-y . hal-01182830

HAL Id: hal-01182830

<https://hal.science/hal-01182830>

Submitted on 3 Aug 2015

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ON THE ELECTROSTATIC BORN-INFELD EQUATION WITH EXTENDED CHARGES

DENIS BONHEURE, PIETRO D'AVENIA, AND ALESSIO POMPONIO

ABSTRACT. In this paper, we deal with the electrostatic Born-Infeld equation

$$(BI) \quad \begin{cases} -\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi|^2}} \right) = \rho, & x \in \mathbb{R}^N, \\ \lim_{|x| \rightarrow \infty} \phi(x) = 0, \end{cases}$$

where ρ is an assigned extended charge density. We are interested in the existence and uniqueness of the potential ϕ and finiteness of the energy of the electrostatic field $-\nabla \phi$. We first relax the problem and treat it with the direct method of the Calculus of Variations for a broad class of charge densities. Assuming ρ is radially distributed, we recover the weak formulation of (BI) and the regularity of the solution of the Poisson equation (under the same smoothness assumptions). In the case of a locally bounded charge, we also recover the weak formulation without assuming any symmetry. The solution is even classical if ρ is smooth. Then we analyze the case where the density ρ is a superposition of point charges and discuss the results in [18]. Other models are discussed, as for instance a system arising from the coupling of the nonlinear Klein-Gordon equation with the Born-Infeld theory.

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2010 *Mathematics Subject Classification.* 35J93, 35Q60, 78A30.

Key words and phrases. Born-Infeld equation, nonlinear electromagnetism, extended charges, mean curvature operator in the Lorentz-Minkowski space.

D. Bonheure is supported by INRIA - Team MEPHYSTO, MIS F.4508.14 (FNRS), PDR T.1110.14F (FNRS) & ARC AUWB-2012-12/17-ULB1- IAPAS. P. d'Avenia and A. Pomponio are supported by GNAMPA Project "Analisi variazionale di modelli fisici non lineari".

1. INTRODUCTION

Classically, the relation between matter and electromagnetic field in the space-time $\mathbb{R} \times \mathbb{R}^3$ can be interpreted from two opposite points of view. Following an unitarian approach, we can consider the electromagnetic field as the unique physical entity and the matter is then given by the singularities of the field. Conversely, according to a dualistic standpoint, field and particle are two different entities: even if the particles generate the field and interact with it, they are not a part of the field. In the dualistic approach, however, according to the original Maxwell theory of electromagnetism, there is the so called infinity problem associated with a point charge source [8, 13]. More precisely, choosing suitably the physical constants, the classical Maxwell equations for the electrostatic case in the vacuum are

$$(1.1) \quad \nabla \times \mathbf{E} = 0,$$

$$(1.2) \quad \operatorname{div} \mathbf{E} = \rho,$$

where \mathbf{E} is the electric field and ρ is the charge density. Since Equation (1.1) implies $\mathbf{E} = -\nabla\phi$, Equation (1.2) yields the so-called Poisson equation

$$(1.3) \quad -\Delta\phi = \rho.$$

Therefore, if $\rho = \delta$, the solution of (1.3) is $\phi(x) = 1/(4\pi|x|)$, but its energy is

$$\mathcal{H} = \frac{1}{2} \int_{\mathbb{R}^3} |\mathbf{E}|^2 dx = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla\phi|^2 dx = +\infty.$$

When $\rho \in L^1(\mathbb{R}^3)$, which is another relevant physical case, we cannot say, in general, that (1.3) admits a solution with finite energy (see e.g. [15] for a counterexample). In fact, it is easily seen from Gagliardo-Nirenberg-Sobolev inequality, see e.g. [19], that the mathematical assumption which implies the finiteness of the energy is $\rho \in L^{6/5}(\mathbb{R}^3)$. This hypothesis does not cover all relevant physical cases.

Maxwell's equations are variational by which we mean that they can be derived as the Euler equations of a Lagrangian. To avoid the violation of the *principle of finiteness*, Max Born proposed a nonlinear theory [5, 6] starting from a modification of Maxwell's Lagrangian density. This theory is built on in analogy with Einstein's mechanics of special relativity. Indeed, one pass from Newton's mechanics to Einstein's mechanics by replacing the action function $\mathcal{L}_N = \frac{1}{2}mv^2$ with $\mathcal{L}_E = mc^2(1 - \sqrt{1 - v^2/c^2})$ as this last expression is one of the simplest which is real only when $v < c$ and gives the classical formulation in the limit of small velocities. By analogy, starting from Maxwell's Lagrangian density in the vacuum

$$(1.4) \quad \mathcal{L}_M = -\frac{F_{\mu\nu}F^{\mu\nu}}{4},$$

where $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$, $(A_0, A_1, A_2, A_3) = (\phi, -\mathbf{A})$ is the electromagnetic potential, $(x_0, x_1, x_2, x_3) = (t, x)$ and ∂_j denotes the partial derivative with respect to x_j , Born introduced the new Lagrangian density

$$(1.5) \quad \mathcal{L}_B = b^2 \left(1 - \sqrt{1 + \frac{F_{\mu\nu}F^{\mu\nu}}{2b^2}} \right) \sqrt{-\det(g_{\mu\nu})},$$

where b is a constant having the dimensions of e/r_0^2 , e and r_0 being respectively the charge and the radius of the electron. In this last formula, $g_{\mu\nu}$ is the Minkowski metric tensor with signature $(+---)$. Since Born's action, as well as Maxwell's action, is invariant only for the Lorentz group of transformations (orthogonal transformations), some months later, Born and Infeld introduced

a modified version of the Lagrangian density

$$(1.6) \quad \mathcal{L}_{\text{BI}} = b^2 \left(\sqrt{-\det(g_{\mu\nu})} - \sqrt{-\det\left(g_{\mu\nu} + \frac{F_{\mu\nu}}{b}\right)} \right),$$

whose integral is now invariant for general transformation [7, 8]. Since the electromagnetic field (\mathbf{E}, \mathbf{B}) is given by

$$\mathbf{B} = \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -\nabla\phi - \partial_t \mathbf{A},$$

(1.4), (1.5) and (1.6) can be written respectively as

$$\mathcal{L}_{\text{M}} = \frac{|\mathbf{E}|^2 - |\mathbf{B}|^2}{2}, \quad \mathcal{L}_{\text{B}} = b^2 \left(1 - \sqrt{1 - \frac{|\mathbf{E}|^2 - |\mathbf{B}|^2}{b^2}} \right)$$

and

$$\mathcal{L}_{\text{BI}} = b^2 \left(1 - \sqrt{1 - \frac{|\mathbf{E}|^2 - |\mathbf{B}|^2}{b^2} - \frac{(\mathbf{E} \cdot \mathbf{B})^2}{b^4}} \right).$$

In the electrostatic case, in which we are interested in this paper, we infer that

$$\mathcal{L}_{\text{B}} = \mathcal{L}_{\text{BI}} = b^2 \left(1 - \sqrt{1 - \frac{|\mathbf{E}|^2}{b^2}} \right).$$

As emphasized above, we recover Newton's classical mechanics from Einstein's special relativity for small velocities or when $c \rightarrow +\infty$. The same holds true with Born-Infeld formulation of electromagnetism: if $b \rightarrow +\infty$ or for electromagnetic fields having small intensities, both \mathcal{L}_{B} and \mathcal{L}_{BI} reduce to Maxwell's Lagrangian density \mathcal{L}_{M} .

In presence of a charge density ρ , we formally get the equation

$$(1.7) \quad -\operatorname{div} \left(\frac{\nabla\phi}{\sqrt{1 - |\nabla\phi|^2/b^2}} \right) = \rho,$$

which replaces the Poisson equation (1.3). This equation can also be obtained observing that the Born-Infeld theory distinguishes between the electric field \mathbf{E} and the electric induction \mathbf{D} : the field \mathbf{D} satisfies

$$(1.8) \quad \operatorname{div} \mathbf{D} = \rho$$

and the fields \mathbf{E} and \mathbf{D} are related by

$$(1.9) \quad \mathbf{D} = \frac{\mathbf{E}}{\sqrt{1 - (|\mathbf{E}|/b)^2}}.$$

Substituting (1.9) in (1.8), we recover (1.7). Finite energy point particle solutions with δ -function sources are known as *BIONS* (see for example [16]). When $\rho = \delta$, one can easily explicitly compute the solution, see for example [22].

From now on, for simplicity and without loss of generality, we fix $b = 1$. It is worth mentioning that the operator Q^- , defined as

$$(1.10) \quad Q^-(\phi) = -\operatorname{div} \left(\frac{\nabla\phi}{\sqrt{1 - |\nabla\phi|^2}} \right),$$

also naturally appears in string theory, in particular in the study of D -branes (see e.g. [16]) and, in classical relativity, where Q^- represents the mean curvature operator in Lorentz-Minkowski space, see for instance [3, 9]. In this last context, the following definition is standard.

Definition 1.1. Let $\phi \in C^{0,1}(\Omega)$, with $\Omega \subset \mathbb{R}^N$. We say that ϕ is

- weakly spacelike if $|\nabla\phi| \leq 1$ a.e. in Ω ;
- spacelike $|\phi(x) - \phi(y)| < |x - y|$ whenever $x, y \in \Omega$, $x \neq y$ and the line segment $\overline{xy} \subset \Omega$;
- strictly spacelike if ϕ is spacelike, $\phi \in C^1(\Omega)$ and $|\nabla\phi| < 1$ in Ω .

Our motivation in this paper is to study rigorously the boundary value problem

$$(BT) \quad \begin{cases} -\operatorname{div} \left(\frac{\nabla\phi}{\sqrt{1-|\nabla\phi|^2}} \right) = \rho, & x \in \mathbb{R}^N, \\ \lim_{|x| \rightarrow \infty} \phi(x) = 0, \end{cases}$$

for general non-trivial charge distributions. Assuming $N \geq 3$, we work in the functional space

$$(1.11) \quad \mathcal{X} = D^{1,2}(\mathbb{R}^N) \cap \{\phi \in C^{0,1}(\mathbb{R}^N) \mid \|\nabla\phi\|_\infty \leq 1\},$$

equipped with the norm defined by

$$\|\phi\|_{\mathcal{X}} := \left(\int_{\mathbb{R}^N} |\nabla\phi|^2 dx \right)^{1/2}.$$

More properties of this space are given in Section 2.1. We recall that $D^{1,2}(\mathbb{R}^N)$ is the completion of $C_c^\infty(\mathbb{R}^N)$ with respect to above norm and we anticipate that \mathcal{X}^* , the dual space of \mathcal{X} , contains Radon measures as for instance superpositions of point charges or $L^1(\mathbb{R}^N)$ densities.

For a $\rho \in \mathcal{X}^*$, *weak solutions* are understood in the following sense.

Definition 1.2. A weak solution of (BT) is a function $\phi_\rho \in \mathcal{X}$ such that for all $\psi \in \mathcal{X}$, we have

$$(1.12) \quad \int_{\mathbb{R}^N} \frac{\nabla\phi_\rho \cdot \nabla\psi}{\sqrt{1-|\nabla\phi_\rho|^2}} dx = \langle \rho, \psi \rangle,$$

where $\langle \cdot, \cdot \rangle$ denotes the duality pairing between \mathcal{X}^* and \mathcal{X} .

Observe that the boundary condition at infinity is encoded in the functional space. We also emphasize that if ρ is a distribution, the weak formulation of (1.12) extends to any test function $\psi \in C_c^\infty(\mathbb{R}^N)$.

As Born-Infeld equation is formally the Euler equation of the action functional $I : \mathcal{X} \rightarrow \mathbb{R}$ defined by

$$(1.13) \quad I(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1-|\nabla\phi|^2} \right) dx - \langle \rho, \phi \rangle,$$

we expect that one can derive existence and uniqueness of the solution from a variational principle. Furthermore, since I is bounded from below in \mathcal{X} and strictly convex, one can look for the solution as the minimizer of I in \mathcal{X} by the direct methods of the Calculus of Variations. However, one needs to pay attention to the lack of regularity of the functional when $\|\nabla\phi\|_\infty = 1$. Hence, as in convex optimization or in the Calculus of Variations for non smooth functionals (see for example [12, 24]), it is natural to relax the notion of critical point: more precisely we say that $\phi_\rho \in \mathcal{X}$ is a *critical point in weak sense* for the functional I if 0 belongs to the subdifferential of I at ϕ_ρ (see Definition 2.4), which, in our case, simply amounts to ask that ϕ_ρ is a minimum for the functional I , see Remark 2.5. We refer to Subsection 2.1 for more details.

We first prove (see Subsection 2.1) existence and uniqueness for the relaxed problem.

Theorem 1.3. For any $\rho \in \mathcal{X}^*$, there exists a unique ϕ_ρ which minimizes I . This is the unique critical point in weak sense of the functional I .

Up to our knowledge, it is not known in the literature whether the weak formulation (1.12) holds or not for critical point in weak sense under the mere assumption $\rho \in \mathcal{X}^*$. This question has motivated several publications in the past years. In [15], the authors deal with the second order expansion of the non smooth part of the functional, assuming $\rho \in L^1(\mathbb{R}^N)$, see Subsection 5.1. In [18], the author considers the special case $\rho = 4\pi \sum_{i=1}^k \alpha_i \delta_{x_i}$, however there is a gap in the proof, see Section 4 for more details.

The first case in which we can deduce the weak formulation is when $\rho \in \mathcal{X}^*$ is a radially distributed charge, see Section 3 for the precise statement.

Theorem 1.4. *If $\rho \in \mathcal{X}^*$ is radially distributed, then there exists a unique (radial) weak solution $\phi_\rho \in \mathcal{X}$ of (BT).*

Under stronger assumptions on ρ , still assuming radial symmetry of the source, we investigate the regularity of the solution and we partially recover the regularity of Poisson equation, for which we refer to [19], (see Theorem 3.2).

In the same Section 3, without symmetry assumptions, we consider the case of locally bounded source.

Theorem 1.5. *If $\rho \in L_{\text{loc}}^\infty(\mathbb{R}^N) \cap \mathcal{X}^*$, then ϕ_ρ is a (locally strictly) space-like weak solution of (BT).*

The case of a superposition of charges, namely

$$\rho = \sum_{i=1}^k a_i \delta_{x_i},$$

where $a_i \in \mathbb{R}$ and $x_i \in \mathbb{R}^N$, for $i = 1, \dots, k$, $k \in \mathbb{N}_0$, is studied in Section 4. We first identify, as in [18], the possible singular points of the solutions. Basically, the minimum ϕ_ρ of I is always strictly spacelike on $\mathbb{R}^N \setminus \Gamma$, where

$$\Gamma = \bigcup_{i \neq j} \overline{x_i x_j}.$$

We then prove that the minimizer is a distributional solution away from the charges and if the intensities are small or if the charges are sufficiently far away from each other, then ϕ_ρ is strictly spacelike on $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$ and singular only and exactly at the location of the charges, i.e.

$$\lim_{x \rightarrow x_i} |\nabla \phi_\rho(x)| = 1.$$

Theorem 1.6. *Assume $\rho = \sum_{i=1}^k a_i \delta_{x_i}$, where $a_i \in \mathbb{R}$ and $x_i \in \mathbb{R}^N$, for all $i = 1, \dots, k$. Then ϕ_ρ is a distributional solution of the Euler-Lagrange equation in $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$. Namely, for every $\psi \in C_c^\infty(\mathbb{R}^N \setminus \{x_1, \dots, x_k\})$, we have*

$$\int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx = \langle \rho, \psi \rangle.$$

It is a classical solution of the equation in $\mathbb{R}^N \setminus \Gamma$, namely $\phi_\rho \in C^\infty(\mathbb{R}^N \setminus \Gamma)$ and

$$-\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi|^2}} \right) = 0$$

in the classical sense in $\mathbb{R}^N \setminus \Gamma$. Moreover,

(1) for any fixed $x_i \in \mathbb{R}^N$, $i = 1, \dots, k$, there exists $\sigma = \sigma(x_1, \dots, x_k) > 0$ such that if

$$\max_{i=1, \dots, k} |a_i| < \sigma,$$

then ϕ_ρ is a classical solution in $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$;

(2) for any $a_i \in \mathbb{R}$, $i = 1, \dots, k$, there exists $\tau = \tau(a_1, \dots, a_k) > 0$ such that if

$$\min_{i, j=1, \dots, k, i \neq j} |x_i - x_j| > \tau,$$

then ϕ_ρ is a classical solution in $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$.

In these last cases, $\phi_\rho \in C^\infty(\mathbb{R}^N \setminus \{x_1, \dots, x_k\})$, it is strictly spacelike on $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$ and

$$\lim_{x \rightarrow x_i} |\nabla \phi_\rho(x)| = 1.$$

Moreover, in some cases, even if we do not know that minimum ϕ_ρ of the functional I is actually a weak solution of (BT), we can say that it is the *limit* of solutions of approximated problems obtained by modifying the differential operator or mollifying the charge density. This is the concern of Section 5.

We conclude in Section 6 by additional results and comments. In particular, we mention how our method completes some previous studies [10, 21, 26, 27] of a field, governed by the nonlinear Klein-Gordon equation, coupled with the electromagnetic field whose Lagrangian density is given by (1.5) or (1.6), by means of the Weil covariant derivatives.

We finally mention that the operator Q^- has been studied in other situations by many authors in the recent years. We refer to [2, 4] for some results in \mathbb{R}^N and to [20] which provides further references for boundary value problems in a bounded domain. In particular Bartnik and Simon have been among the first to deal with this type of differential operator and some of the ideas from [3] are fundamental in our arguments.

We close the introduction with some notations: C denotes a generic positive constants which can change from line to line, B_R is the ball centered in 0 with radius $R > 0$ and ω_N denotes the measure of the $(N - 1)$ -dimensional unitarian sphere. For every $1 \leq p < N$, p^* is the critical Sobolev exponent in the Sobolev inequality, namely $p^* = Np/(N - p)$.

Acknowledgement. This work has been partially carried out during a stay of P.D. and A.P. in Bruxelles. They would like to express their deep gratitude to the Département de Mathématique, Université libre de Bruxelles, for the support and warm hospitality. D.B. acknowledge the support of INDAM for his visits at the Politecnico di Bari where parts of this work have been achieved.

2. A RELAXED FORMULATION VIA NONSMOOTH ANALYSIS

This section is devoted to the proof of Theorem 1.3 and to useful properties of the minimum ϕ_ρ . We start with the functional setting and we recall some well-known facts from convex analysis and non smooth critical point theory.

2.1. Functional setting and the existence of the critical point in weak sense. We start with some properties of the ambient space \mathcal{X} defined in (1.11). The proof follows from standard arguments that we give for completeness.

Lemma 2.1. *The following assertions hold:*

- (i) \mathcal{X} is continuously embedded in $W^{1,p}(\mathbb{R}^N)$, for all $p \geq 2^* = 2N/(N - 2)$;
- (ii) \mathcal{X} is continuously embedded in $L^\infty(\mathbb{R}^N)$;
- (iii) if $\phi \in \mathcal{X}$, then $\lim_{|x| \rightarrow \infty} \phi(x) = 0$;
- (iv) \mathcal{X} is weakly closed;

(v) if $(\phi_n)_n \subset \mathcal{X}$ is bounded, there exists $\bar{\phi} \in \mathcal{X}$ such that, up to a subsequence, $\phi_n \rightharpoonup \bar{\phi}$ weakly in \mathcal{X} and uniformly on compact sets.

Proof. By definition, if $\phi \in \mathcal{X}$, then for every $q \geq 2$, we have $|\nabla \phi| \in L^q(\mathbb{R}^N)$ whence $\phi \in L^{q^*}(\mathbb{R}^N)$ for $q \in [2, N)$ by Sobolev inequality. Hence $\phi \in W^{1,p}(\mathbb{R}^N)$ for $p \geq 2^*$. The continuity of the imbedding is clear and then the proof of assertion (i) is complete.

Assertions (ii) and (iii) are direct consequences of (i), Morrey-Sobolev inequality.

As regards (iv), since \mathcal{X} is convex, it is sufficient to show that \mathcal{X} is closed with respect to the strong topology. Take $(\phi_n)_n \subset \mathcal{X}$ such that $\phi_n \rightarrow \phi$ in \mathcal{X} . Then we have

$$|\phi_n(x) - \phi_n(y)| \leq |x - y|,$$

for all $x, y \in \mathbb{R}^N$. Since $\phi_n \rightarrow \phi$ uniformly in \mathbb{R}^N by (ii), we conclude that $\|\nabla \phi\|_\infty \leq 1$. We finally prove (v). Since $(\phi_n)_n$ is a bounded sequence in $D^{1,2}(\mathbb{R}^N)$, it contains a weakly converging subsequence that we still denote by $(\phi_n)_n$. By (iv), the weak limit $\bar{\phi}$ belongs to \mathcal{X} and by Ascoli-Arzelà Theorem, the convergence is uniform on compact sets. \square

We now give some properties of the functional I . The simple inequality

$$(2.1) \quad \frac{1}{2}t \leq 1 - \sqrt{1-t} \leq t, \quad \text{for all } t \in [0, 1]$$

will be useful.

Lemma 2.2. *The functional $I : \mathcal{X} \rightarrow \mathbb{R}$ is*

- (i) *bounded from below,*
- (ii) *coercive,*
- (iii) *continuous,*
- (iv) *strictly convex,*
- (v) *weakly lower semi-continuous.*

Proof. Using (2.1), we infer that

$$I(\phi) \geq \frac{1}{2}\|\nabla \phi\|_2^2 - \|\rho\|_{\mathcal{X}^*}\|\nabla \phi\|_2$$

for every $\phi \in \mathcal{X}$ and this yields (i) and (ii). As regards (iii), we only need to prove that $J : \mathcal{X} \rightarrow \mathbb{R}$, defined by

$$J(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx,$$

is continuous. To this end, we consider a sequence $(\phi_n)_n \subset \mathcal{X}$ that converges to ϕ in \mathcal{X} . Then, up to a subsequence, $\nabla \phi_n \rightarrow \nabla \phi$ a.e. in \mathbb{R}^N and there exists $w \in L^1(\mathbb{R}^N)$ such that $|\nabla \phi_n|^2 \leq w$, $|\nabla \phi|^2 \leq w$ a.e. in \mathbb{R}^N . Thus, by (2.1),

$$\left| (1 - \sqrt{1 - |\nabla \phi_n|^2}) - (1 - \sqrt{1 - |\nabla \phi|^2}) \right| \leq |\nabla \phi_n|^2 + |\nabla \phi|^2 \leq 2w$$

and then, by Lebesgue's Dominated Convergence Theorem we have that $J(\phi_n) \rightarrow J(\phi)$. It is straightforward to check that the convergence holds for the whole sequence.

The strict convexity of J follows from the strict convexity of the real function $y \in B_1 \mapsto 1 - \sqrt{1 - |y|^2}$. Since J is continuous and convex whereas ρ is continuous with respect to the weak convergence, the weak lower semi-continuity holds. \square

As a consequence of the previous properties, we get the existence of a unique minimizer.

Proposition 2.3. *The infimum $m = \inf_{\phi \in \mathcal{X}} I(\phi)$ is achieved by a unique $\phi_\rho \in \mathcal{X} \setminus \{0\}$.*

Proof. The existence and uniqueness follow from Lemma 2.2. Therefore, we only have to show that ϕ_ρ is nontrivial or equivalently that $m < 0$. Taking $\phi \in \mathcal{X}$ such that $\langle \rho, \phi \rangle > 0$, we compute

$$I(t\phi) \leq t^2 \|\nabla \phi\|_2^2 - t \langle \rho, \phi \rangle < 0,$$

for $t > 0$ small enough, whence $m < 0$. \square

We now recall some classical definitions from convex analysis, see [24].

Definition 2.4. Let X be a real Banach space and $\psi : X \rightarrow (-\infty, +\infty]$ be a convex lower semicontinuous function. Let $D(\psi) = \{u \in X \mid \psi(u) < +\infty\}$ be the effective domain of ψ . For $u \in D(\psi)$, the set

$$\partial\psi(u) = \{u^* \in X^* \mid \psi(v) - \psi(u) \geq \langle u^*, v - u \rangle, \forall v \in X\}$$

is called the subdifferential of ψ at u . If, moreover, we consider a functional $I = \psi + \Phi$, with ψ as above and $\Phi \in C^1(X, \mathbb{R})$, then $u \in D(\psi)$ is said to be critical in weak sense if $-\Phi'(u) \in \partial\psi(u)$, that is

$$\langle \Phi'(u), v - u \rangle + \psi(v) - \psi(u) \geq 0, \forall v \in X.$$

Remark 2.5. Observe that, according to the previous definition, ϕ_ρ is a critical point in weak sense for the functional I if and only if, for any $\phi \in \mathcal{X}$ we get

$$\int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx - \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi_\rho|^2}\right) dx \geq \langle \rho, \phi - \phi_\rho \rangle,$$

which is simply equivalent to require that ϕ_ρ is a minimum for I .

Proposition 2.3 and Remark 2.5 leads easily to Theorem 1.3.

2.2. Further properties. Here we present some properties that are useful in the sequel. We start by recalling the convexity of the functional implies that weak solutions are minimizers and therefore we deduce the uniqueness of the weak solution (if any). As before, for $\rho \in \mathcal{X}^*$, we denote the unique minimizer of I in \mathcal{X} by ϕ_ρ .

Proposition 2.6. Assume $\rho \in \mathcal{X}^*$. If $\phi \in \mathcal{X}$ is a weak solution of (BT), then $\phi = \phi_\rho$.

Proof. If $\phi \in \mathcal{X}$ satisfies (1.12), it is easy to see that $|\nabla \phi(x)| < 1$ for a.e. $x \in \mathbb{R}^N$. By convexity, we get

$$(2.2) \quad 1 - \sqrt{1 - |\nabla \phi_\rho(x)|^2} \geq 1 - \sqrt{1 - |\nabla \phi(x)|^2} + \frac{\nabla \phi(x) \cdot (\nabla \phi_\rho(x) - \nabla \phi(x))}{\sqrt{1 - |\nabla \phi(x)|^2}}$$

for a.e. $x \in \mathbb{R}^N$. Moreover, again since ϕ satisfies (1.12), we have

$$(2.3) \quad \int_{\mathbb{R}^N} \frac{|\nabla \phi|^2}{\sqrt{1 - |\nabla \phi|^2}} dx - \int_{\mathbb{R}^N} \frac{\nabla \phi \cdot \nabla \phi_\rho}{\sqrt{1 - |\nabla \phi|^2}} dx = \langle \rho, \phi - \phi_\rho \rangle.$$

Combining (2.2) and (2.3), we conclude that $I(\phi_\rho) \geq I(\phi)$. Uniqueness of the minimizer ϕ_ρ of I in \mathcal{X} leads to the conclusion. \square

Proposition 2.6 relies on the fact that weak solutions are minimizers. A question now arises naturally: does the reverse statement holds? Namely, is it true that the unique minimizer ϕ_ρ is always a weak solution of (BT)? We are not able to answer this question in its full generality but we conjecture a positive answer and the following statement goes in that direction.

Proposition 2.7. *Assume $\rho \in \mathcal{X}^*$ and let ϕ_ρ be the unique minimizer of I in \mathcal{X} . Then*

$$E = \{x \in \mathbb{R}^N \mid |\nabla \phi_\rho| = 1\}$$

is a null set (with respect to Lebesgue measure) and the function ϕ_ρ satisfies

$$(2.4) \quad \int_{\mathbb{R}^N} \frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx \leq \langle \rho, \phi_\rho \rangle.$$

Moreover, for all $\psi \in \mathcal{X}$, we have the variational inequality

$$(2.5) \quad \int_{\mathbb{R}^N} \frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx - \int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx \leq \langle \rho, \phi_\rho - \psi \rangle.$$

Proof. Since for every $t \in [0, 1]$ and $\psi \in \mathcal{X}$, $\phi_t = \phi_\rho + t(\psi - \phi_\rho) \in \mathcal{X}$, we have $I(\phi_\rho) \leq I(\phi_t)$, namely

$$(2.6) \quad \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi_\rho(x)|^2}\right) dx - \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi_t(x)|^2}\right) dx \leq t \langle \rho, \phi_\rho - \psi \rangle.$$

In the particular case $\psi = 0$, we have

$$\int_{E^c} \frac{\sqrt{1 - (1-t)^2 |\nabla \phi_\rho|^2} - \sqrt{1 - |\nabla \phi_\rho|^2}}{t} dx + \frac{\sqrt{2-t}}{\sqrt{t}} |E| \leq \langle \rho, \phi_\rho \rangle$$

for all $t \in (0, 1]$. Since both terms in the left hand side are nonnegative, we infer that $|E| = 0$ and moreover

$$\int_{\mathbb{R}^N} \frac{(2-t)|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2} + \sqrt{1 - (1-t)^2 |\nabla \phi_\rho|^2}} dx \leq \langle \rho, \phi_\rho \rangle.$$

By Fatou's Lemma, letting t go to zero, we get (2.4) and therefore

$$(2.7) \quad \frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} \in L^1(\mathbb{R}^N).$$

Considering now $\psi \neq 0$ in (2.6), as $|E| = 0$, we deduce that

$$(2.8) \quad \int_{\mathbb{R}^N} \frac{(2-t)|\nabla \phi_\rho|^2 - 2(1-t)\nabla \phi_\rho \cdot \nabla \psi - t|\nabla \psi|^2}{\sqrt{1 - |\nabla \phi_\rho|^2} + \sqrt{1 - |\nabla \phi_t|^2}} dx \leq \langle \rho, \phi_\rho - \psi \rangle.$$

Aiming to apply Lebesgue's Dominated Convergence Theorem in (2.8), we first notice that

$$\left| \frac{(2-t)|\nabla \phi_\rho|^2 - 2(1-t)\nabla \phi_\rho \cdot \nabla \psi - t|\nabla \psi|^2}{\sqrt{1 - |\nabla \phi_\rho|^2} + \sqrt{1 - |\nabla \phi_t|^2}} \right| \leq C \left(\frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} + \frac{|\nabla \psi|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} \right)$$

The first term of the right hand side is $L^1(\mathbb{R}^N)$ by (2.7). To estimate the second one, we observe that

$$\int_{A_\rho} \frac{|\nabla \psi|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx \leq C \int_{\mathbb{R}^N} |\nabla \psi|^2 dx$$

and

$$\int_{A_\rho^c} \frac{|\nabla \psi|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx \leq C \int_{\mathbb{R}^N} \frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx,$$

where we have set $A_\rho := \{x \in \mathbb{R}^N \mid |\nabla \phi_\rho| \leq 1/2\}$. Since we are now allowed to take the limit as $t \rightarrow 0^+$ in (2.8), we get (2.5). \square

Remark 2.8. If ϕ_ρ satisfies further

$$\int_{\mathbb{R}^N} \frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx = \langle \rho, \phi_\rho \rangle,$$

then, by (2.5), it is easy to see that ϕ_ρ is a weak solution of (BT).

Remark 2.9. It is clear from the proof of Proposition 2.7 that it is enough to assume $\psi \in D^{1,2}(\mathbb{R}^N)$ together with $|\nabla \psi| \in L^\infty(\mathbb{R}^N)$ to get

$$\int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx \in \mathbb{R}.$$

The next lemma states a useful convergence to prove Theorem 1.4.

Lemma 2.10. Assume $\rho \in \mathcal{X}^*$ and let ϕ_ρ be the unique minimizer of I in \mathcal{X} . If $(\psi_n)_n \subset D^{1,2}(\mathbb{R}^N)$ is such that $\|\nabla \psi_n\|_\infty \leq C$ for some $C > 0$ and $\psi_n \rightarrow \psi$ in $D^{1,2}(\mathbb{R}^N)$ then, up to a subsequence,

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi_n}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx = \int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx.$$

Proof. Keeping the notation $A_\rho = \{x \in \mathbb{R}^N \mid |\nabla \phi_\rho| \leq 1/2\}$, we have

$$\int_{A_\rho} \frac{\nabla \phi_\rho \cdot (\nabla \psi_n - \nabla \psi)}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx \leq C \left(\int_{\mathbb{R}^N} |\nabla \phi_\rho|^2 dx \right)^{1/2} \left(\int_{\mathbb{R}^N} |\nabla(\psi_n - \psi)|^2 dx \right)^{1/2}$$

whereas, on A_ρ^c ,

$$\left| \frac{\nabla \phi_\rho \cdot (\nabla \psi_n - \nabla \psi)}{\sqrt{1 - |\nabla \phi_\rho|^2}} \right| \leq C \frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}}.$$

Recalling that

$$\frac{|\nabla \phi_\rho|^2}{\sqrt{1 - |\nabla \phi_\rho|^2}} \in L^1(\mathbb{R}^N),$$

we can apply Lebesgue's Dominated Convergence Theorem on A_ρ^c so that the conclusion follows. \square

In the spirit of [3, Lemma 1.2], we next provide a monotonicity property of the map $\rho \in \mathcal{X}^* \mapsto \phi_\rho$, where ϕ_ρ is the unique minimum associated to ρ . We first order the elements of \mathcal{X}^* .

Definition 2.11. Let $\rho_1, \rho_2 \in \mathcal{X}^*$. We say that $\rho_1 \leq \rho_2$, if we have $\langle \rho_1, \varphi \rangle \leq \langle \rho_2, \varphi \rangle$ for any $\varphi \in \mathcal{X}$ with $\varphi \geq 0$.

The next lemma is a *comparison principle* for minimizers.

Lemma 2.12. If $\rho_1, \rho_2 \in \mathcal{X}^*$ are such that $\rho_1 \leq \rho_2$, then $\phi_{\rho_1}(x) \leq \phi_{\rho_2}(x)$ for all $x \in \mathbb{R}^N$.

Proof. Let $\phi_i = \phi_{\rho_i}$. Suppose by contradiction that the open set $\Omega^+ := \{x \in \mathbb{R}^N \mid \phi_1(x) > \phi_2(x)\}$ is non-empty. Of course, by continuity, we have $\phi_1 = \phi_2$ on $\partial\Omega^+$. We define $\Omega^- := \mathbb{R}^N \setminus \Omega^+$, which may be the empty set and we introduce the functionals $I_1, I_2 : \mathcal{X} \rightarrow \mathbb{R}$, where

$$I_i(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2} \right) dx - \langle \rho_i, \phi \rangle.$$

We also fix the notation

$$J_\pm(\phi) = \int_{\Omega^\pm} \left(1 - \sqrt{1 - |\nabla \phi|^2} \right) dx.$$

Let

$$\psi_1(x) = \begin{cases} \phi_2(x) & \text{in } \Omega^+, \\ \phi_1(x) & \text{in } \Omega^-, \end{cases}$$

and

$$\psi_2(x) = \begin{cases} \phi_1(x) & \text{in } \Omega^+, \\ \phi_2(x) & \text{in } \Omega^-. \end{cases}$$

Clearly, $\psi_i \in \mathcal{X}$ and $\psi_2 - \phi_2 = \phi_1 - \psi_1 \geq 0$. Then

$$\begin{aligned} I_2(\psi_2) &= J_+(\phi_1) + J_-(\phi_2) - \langle \rho_2, \psi_2 \rangle \\ &= I_2(\phi_2) + J_+(\phi_1) - J_+(\phi_2) - \langle \rho_2, \phi_1 - \psi_1 \rangle \\ &\leq I_2(\phi_2) + J_+(\phi_1) - J_+(\phi_2) - \langle \rho_1, \phi_1 - \psi_1 \rangle \\ &= I_2(\phi_2) + I_1(\phi_1) - I_1(\psi_1). \end{aligned}$$

Now, since ϕ_1 is the unique minimizer of I_1 and ψ_1 does not coincide with ϕ_1 on an open set, we conclude that

$$I_2(\phi_2) + I_1(\phi_1) - I_1(\psi_1) < I_2(\phi_2),$$

and we reach a contradiction with the minimality of ϕ_2 . \square

We conclude this section with the following property which will be useful in the sequel.

Lemma 2.13. *Let $(\rho_n)_n \subset \mathcal{X}^*$ be a bounded sequence. Then there exists $\bar{\phi} \in \mathcal{X}$ such that $\phi_{\rho_n} \rightharpoonup \bar{\phi}$ weakly in \mathcal{X} and uniformly on compact sets.*

Proof. For short, we set $\phi_n = \phi_{\rho_n}$ and

$$I_n(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx - \langle \rho_n, \phi \rangle.$$

By Assertion (v) of Lemma 2.1, it is enough to show that $(\phi_n)_n$ is bounded. Since $(\rho_n)_n$ is a bounded sequence in \mathcal{X}^* and $I_n(\phi_n) \leq 0$, we have

$$\frac{1}{2} \|\nabla \phi_n\|_2^2 \leq \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi_n|^2}\right) dx \leq \langle \rho_n, \phi_n \rangle \leq C \|\nabla \phi_n\|_2,$$

namely $(\phi_n)_n$ is a bounded sequence in \mathcal{X} . \square

3. WEAK FORMULATION FOR RADIALLY DISTRIBUTED OR BOUNDED CHARGE DENSITIES

3.1. The case of a radial charge. We now turn our attention to radially distributed charges. In this case, we are able to recover the weak formulation from the relaxed equation. Indeed we build a dense set, within the set of test functions, of admissible variations. The argument is borrowed from [23] where it was used to handle a monotonicity constraint.

We first precise the meaning of radially distributed charge density. For $\tau \in O(N)$, $\phi \in \mathcal{X}$ and $\rho \in \mathcal{X}^*$, we define $\phi^\tau \in \mathcal{X}$ as $\phi^\tau(x) = \phi(\tau x)$, for all $x \in \mathbb{R}^N$, and $\rho^\tau \in \mathcal{X}^*$ as $\langle \rho^\tau, \psi \rangle = \langle \rho, \psi^\tau \rangle$, for all $\psi \in \mathcal{X}$.

Definition 3.1. *We say that $\rho \in \mathcal{X}^*$ is radially distributed if $\rho^\tau = \rho$, for any $\tau \in O(N)$.*

We next define

$$\mathcal{X}_{\text{rad}} = \{\phi \in \mathcal{X} \mid \phi^\tau = \phi \text{ for every } \tau \in O(N)\}$$

and if $\phi \in \mathcal{X}_{\text{rad}}$, in order to simplify the notations, we keep ϕ to denote the function $r = |x| \in \mathbb{R}_+ \mapsto \phi(r)$ of a single real variable. We make furthermore a similar identification for the radially distributed maps $\rho \in \mathcal{X}^*$.

Given these definitions, we are ready to prove that radial minimizers are weak solutions.

Proof of Theorem 1.4. As starting point, let us show that $\phi_\rho \in \mathcal{X}_{\text{rad}}$. Indeed, for any $\tau \in O(N)$, since ρ radially distributed, we have $I(\phi_\rho^\tau) = I(\phi_\rho)$ and so we conclude by the uniqueness of the minimum.

We now prove that ϕ_ρ is a weak solution of (BI). Define, for $k \in \mathbb{N} \setminus \{0\}$, the sets

$$E_k = \left\{ r \geq 0 \mid |\phi'_\rho(r)| \geq 1 - \frac{1}{k} \right\}.$$

Since

$$E = \{r \geq 0 \mid |\phi'_\rho(r)| = 1\}$$

is a null set, we have that $|\bigcap_{k \geq 1} E_k| = 0$.

Take $\psi \in \mathcal{X}_{\text{rad}} \cap C_c^\infty(\mathbb{R}^N)$ with $\text{supp } \psi \subset [0, R]$ and let

$$\psi_k(r) = - \int_r^{+\infty} \psi'(s) [1 - \chi_{E_k}(s)] ds.$$

Of course, by construction, $\text{supp } \psi_k \subset [0, R]$, for any $k \geq 1$. Moreover, if $|t|$ is sufficiently small, then $\phi_\rho + t\psi_k \in \mathcal{X}$. Indeed, since $(\phi_\rho + t\psi_k)' = \phi'_\rho + t\psi'[1 - \chi_{E_k}]$, then, if $r \in E_k$

$$|(\phi_\rho + t\psi_k)'(r)| = |\phi'_\rho(r)| \leq 1,$$

otherwise, taking $|t| \leq \frac{1}{k\|\psi'\|_\infty}$, we have

$$|(\phi_\rho + t\psi_k)'(r)| \leq |\phi'_\rho(r)| + |t|\|\psi'\|_\infty < 1 - \frac{1}{k} + |t|\|\psi'\|_\infty \leq 1.$$

Now, since ϕ_ρ is the minimizer of I , arguing as in the proof of Proposition 2.7, we infer that

$$(3.1) \quad \lim_{t \rightarrow 0} \frac{I(\phi_\rho + t\psi_k) - I(\phi_\rho)}{t} = \omega_N \int_0^{+\infty} \frac{\phi'_\rho \psi'}{\sqrt{1 - |\phi'_\rho|^2}} [1 - \chi_{E_k}] r^{N-1} dr - \langle \rho, \psi_k \rangle = 0,$$

for every $k \geq 1$. Since $E_{k+1} \subset E_k$ and $|E_k| \rightarrow 0$, as $k \rightarrow +\infty$, we have that $\chi_{E_k} \rightarrow 0$ a.e. in \mathbb{R}^N and so, by Lebesgue's Dominated Convergence Theorem,

$$\int_0^{+\infty} \frac{\phi'_\rho \psi'}{\sqrt{1 - |\phi'_\rho|^2}} [1 - \chi_{E_k}] r^{N-1} dr \rightarrow \int_0^{+\infty} \frac{\phi'_\rho \psi'}{\sqrt{1 - |\phi'_\rho|^2}} r^{N-1} dr.$$

Moreover, since it is easily seen that $\psi_k \rightarrow \psi$ in \mathcal{X} , we have

$$\langle \rho, \psi_k \rangle \rightarrow \langle \rho, \psi \rangle,$$

as $k \rightarrow \infty$. Hence, for any $\psi \in \mathcal{X}_{\text{rad}} \cap C_c^\infty(\mathbb{R}^N)$, taking the limit in (3.1) as $k \rightarrow \infty$, we conclude that

$$(3.2) \quad \int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx = \langle \rho, \psi \rangle.$$

We finally show that (3.2) holds also for any $\psi \in \mathcal{X}_{\text{rad}}$. Given $\psi \in \mathcal{X}_{\text{rad}}$, it is easy to see that there exists $(\psi_n)_n \subset C_c^\infty(\mathbb{R}^N)$, ψ_n radially symmetric such that $\psi_n \rightarrow \psi$ in $D^{1,2}(\mathbb{R}^N)$ and with $\|\nabla \psi_n\|_\infty \leq C$. Indeed it is sufficient to consider $\psi_n = \zeta_n * (\chi_n \psi)$, where ζ_n are smooth radially symmetric mollifiers with compact support, and $\chi_n = \chi(\cdot/n)$, where $\chi : \mathbb{R}^N \rightarrow \mathbb{R}$ is a smooth radially symmetric function such that

$$\chi(x) = \begin{cases} 1 & \text{if } |x| \leq 1, \\ 0 & \text{if } |x| \geq 2. \end{cases}$$

Then, by (3.2), we have

$$\int_{\mathbb{R}^N} \frac{\nabla \phi_\rho \cdot \nabla \psi_n}{\sqrt{1 - |\nabla \phi_\rho|^2}} dx = \langle \rho, \psi_n \rangle,$$

for any $n \geq 1$, and since $\psi_n \rightarrow \psi$ in $D^{1,2}(\mathbb{R}^N)$, with $\|\nabla \psi_n\|_\infty \leq C$, we infer from Lemma 2.10 that (3.2) holds for any $\psi \in \mathcal{X}_{\text{rad}}$. To conclude, it remains to show that (3.2) holds for any $\psi \in \mathcal{X}$.

Since ϕ_ρ is radially symmetric, we can take $\psi = \phi_\rho$ in (3.2) and therefore Remark 2.8 allows to conclude. \square

Assuming further hypotheses on ρ , we can prove that the solution is C^1 .

Theorem 3.2. *Assume that ρ is a radially symmetric function such that $\rho \in L^s(\mathbb{R}^N) \cap L^\sigma(B_\delta(0))$, for some $s \geq 1$, $\sigma \geq N$ and $\delta > 0$. Then the weak solution ϕ_ρ of (BT) is $C^1(\mathbb{R}^N; \mathbb{R})$.*

We start from the following proposition that basically states that $|\phi'_\rho(r)| = 1$ can only happen at $r = 0$.

Proposition 3.3. *Assume that $\rho \in L^s(\mathbb{R}^N)$ with $s \geq 1$ and ρ is radially symmetric. Then, $\phi_\rho \in C^1(0, +\infty)$. Moreover, for every $0 < r_0 < R$ there exists $\varepsilon > 0$ such that $|\phi'_\rho(r)| \leq 1 - \varepsilon$, for every $r \in [r_0, R]$.*

Proof. From the weak formulation written in radial coordinates, we have

$$(3.3) \quad \int_0^{+\infty} \frac{\phi'_\rho(r) r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} \psi'(r) dr = \int_0^{+\infty} r^{N-1} \rho(r) \psi(r) dr,$$

for every $\psi \in \mathcal{X}_{rad}$. We claim that for any $R > 0$, the function

$$(3.4) \quad h_R(r) := \frac{\phi'_\rho(r) r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}}$$

is continuous on $[0, R]$. From this claim we get the regularity of ϕ_ρ and we deduce that for every $R > 0$, there exists a constant $C = C(R) > 0$ such that, for every $r \in (0, R]$

$$\frac{|\phi'_\rho(r)|}{\sqrt{1 - |\phi'_\rho(r)|^2}} \leq \frac{C(R)}{r^{N-1}}.$$

Consequently, for every $r_0 \in (0, R]$, there exists $\varepsilon > 0$ such that

$$|\phi'_\rho(r)| \leq 1 - \varepsilon,$$

for every $r \in [r_0, R]$.

To prove our claim, observe that as $\rho \in L^s(\mathbb{R}^N)$ for some $s \geq 1$, we deduce from (3.3), that

$$\left(\frac{\phi'_\rho(r) r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} \right)' \in L^1(0, R),$$

for any $R > 0$. Moreover, by (2.7), we have

$$\begin{aligned} \int_0^R \frac{|\phi'_\rho(r)| r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} dr &= \int_{[0, R] \cap \{|\phi'_\rho(r)| \leq \frac{1}{2}\}} \frac{|\phi'_\rho(r)| r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} dr + \int_{[0, R] \cap \{|\phi'_\rho(r)| > \frac{1}{2}\}} \frac{|\phi'_\rho(r)| r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} dr \\ &\leq C \left(\int_0^R r^{N-1} dr + \int_0^R \frac{|\phi'_\rho(r)|^2 r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} dr \right) < +\infty. \end{aligned}$$

It follows that $h_R \in W^{1,1}(0, R)$ and the claim is therefore a consequence of the continuous embedding of $W^{1,1}(0, R)$ in $C([0, R])$. \square

Assuming moreover that $\rho \in L^\sigma(B_\delta(0))$ for some $\sigma \geq N$ and $\delta > 0$, we can improve the regularity up to the origin.

Proof of Theorem 3.2. Let $R > 0$ and $\psi_R : \mathbb{R}^N \rightarrow \mathbb{R}$ be defined as

$$\psi_R(x) = \begin{cases} R - |x| & |x| \leq R, \\ 0 & |x| > R. \end{cases}$$

By (1.12), we get

$$(3.5) \quad -\frac{1}{R} \int_0^R \frac{\phi'_\rho(r) r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} dr = \int_0^R \rho(r) r^{N-1} dr - \frac{1}{R} \int_0^R \rho(r) r^N dr.$$

By the Mean Value Theorem (remember that h_R defined in (3.4) is continuous), we have

$$\lim_{R \rightarrow 0^+} -\frac{1}{R} \int_0^R \frac{\phi'_\rho(r) r^{N-1}}{\sqrt{1 - |\phi'_\rho(r)|^2}} dr = -\lim_{R \rightarrow 0^+} \frac{\phi'_\rho(R) R^{N-1}}{\sqrt{1 - |\phi'_\rho(R)|^2}} = c_0 \in \mathbb{R}.$$

Moreover, if R tends to 0, the right hand side of (3.5) tends to 0 by the absolute continuity of the integral. Therefore $c_0 = 0$.

The equation (3.3) means that the weak derivative of $h_R(r)$ is equal to $-\rho(r)r^{N-1}$. Since ρ is L^1_{loc} , we can integrate h'_R on $[0, R]$ to get

$$\frac{\phi'_\rho(R) R^{N-1}}{\sqrt{1 - |\phi'_\rho(R)|^2}} = -\int_0^R \rho(r) r^{N-1} dr,$$

since $c_0 = 0$. Then, if $R < \delta$,

$$\frac{|\phi'_\rho(R)|}{\sqrt{1 - |\phi'_\rho(R)|^2}} \leq \frac{1}{R^{N-1}} \int_0^R |\rho(r)| r^{N-1} dr \leq C R^{(\sigma-N)/\sigma} \|\rho\|_{L^\sigma(B_R(0))}.$$

Since $\sigma \geq N$, we deduce again from the absolute continuity of the integral that $\phi'_\rho(R) \rightarrow 0$, as $R \rightarrow 0^+$. Hence $\phi_\rho \in C^1([0, \infty))$ with $\phi'_\rho(0) = 0$. We conclude therefore that $\phi_\rho \in C^1(\mathbb{R}^N; \mathbb{R})$. \square

3.2. The case of a bounded charge. Here, keeping in mind Definition 1.1, we prove Theorem 1.5, whose assumption is $\rho \in L^\infty_{\text{loc}}(\mathbb{R}^N) \cap \mathcal{X}^*$. Let Ω be an arbitrary bounded domain with smooth boundary in \mathbb{R}^N . We set

$$C^{0,1}_{\phi_\rho}(\Omega) = \{\phi \in C^{0,1}(\Omega) \mid \phi|_{\partial\Omega} = \phi_\rho|_{\partial\Omega}, |\nabla\phi| \leq 1\},$$

$$(3.6) \quad K = \{\overline{xy} \subset \Omega \mid x, y \in \partial\Omega, x \neq y, |\phi_\rho(x) - \phi_\rho(y)| = |x - y|\},$$

and define $I_\Omega : C^{0,1}_{\phi_\rho}(\Omega) \rightarrow \mathbb{R}$ by

$$I_\Omega(\phi) = \int_\Omega \left(1 - \sqrt{1 - |\nabla\phi|^2}\right) dx - \int_\Omega \rho\phi dx.$$

Proof of Theorem 1.5. Using the above notations, for any fixed Ω , it is easy to see that $\phi_\rho|_{\partial\Omega}$ is a minimizer for I_Ω in $C^{0,1}_{\phi_\rho}(\Omega)$. By [3, Corollary 4.2], we have that ϕ_ρ is strictly spacelike in $\Omega \setminus K$ and $Q^-(\phi_\rho) = \rho$ in $\Omega \setminus K$, where Q^- is defined in (1.10). Furthermore,

$$\phi_\rho(tx + (1-t)y) = t\phi_\rho(x) + (1-t)\phi_\rho(y), \quad 0 < t < 1$$

for every $x, y \in \partial\Omega$ such that $|\phi_\rho(x) - \phi_\rho(y)| = |x - y|$ and $\overline{xy} \subset \Omega$. If $K = \emptyset$, then ϕ_ρ is strictly spacelike in Ω .

Assume by contradiction that $K \neq \emptyset$. Then there exist $x, y \in \partial\Omega$ such that $x \neq y$, $\overline{xy} \subset \Omega$ and

$|\phi_\rho(x) - \phi_\rho(y)| = |x - y|$. Without loss of generality we can assume that $\phi_\rho(x) > \phi_\rho(y)$. It is easy to see that for all $t \in (0, 1)$

$$(3.7) \quad \phi_\rho(tx + (1 - t)y) = \phi_\rho(y) + t|x - y|.$$

Since, for any $R > 0$ such that $\Omega \subset B_R$, $\phi_\rho|_{B_R}$ is a minimizer of I_{B_R} in $C_{\phi_\rho}^{0,1}(B_R)$, then, by [3, Theorem 3.2], we have that (3.7) holds for all $t \in \mathbb{R}$ such that $tx + (1 - t)y \in B_R$. Now we reach a contradiction with the boundedness of ϕ_ρ , for an R sufficiently large. \square

Remark 3.4. As observed in [3, Remark p. 147], if $\rho \in C^k(\mathbb{R}^N)$, then $\phi_\rho \in C^{k+1}(\mathbb{R}^N)$.

4. THE ELECTRIC POTENTIAL PRODUCED BY A DISTRIBUTION OF k POINT CHARGES

This section is devoted to the proof of Theorem 1.6. In all this section we set

$$\rho = \sum_{i=1}^k a_i \delta_{x_i},$$

where $a_i \in \mathbb{R}$ and $x_i \in \mathbb{R}^N$, for $i = 1, \dots, k$, $k \in \mathbb{N}_0$. We consider the problem

$$(4.1) \quad \begin{cases} -\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi|^2}} \right) = \sum_{i=1}^k a_i \delta_{x_i}, & \text{in } \mathbb{R}^N, \\ \phi(x) \rightarrow 0, & \text{as } x \rightarrow \infty. \end{cases}$$

In the recent contribution [18], the author claims the existence [18, Proposition 2.1] of a weak solution v_∞ of (4.1) under the assumption that $x_i \in \mathbb{R}^3$ and $a_i \in \mathbb{R}$ for $i = 1, \dots, k$. However, the proof of [18, Step 2.6, page 515] is incomplete because even if the Lebesgue measure of the set $\{x \in \mathbb{R}^3 \mid |\nabla v_\infty(x)| = 1\}$ is zero, yet one cannot compute the variation $\mathcal{F}^{(1)}[v_\infty](\psi)$ for all test functions ψ . Therefore, [18, (22) page 516] does only vanish for a restricted set of test functions which is not necessarily a dense set of $C_c^\infty(\mathbb{R}^N)$. This means that one cannot conclude that v_∞ weakly solves (4.1). As a consequence, the arguments of [18] based on the weak formulation of (4.1) need to be adjusted and the existence of a weak solution of (4.1) remains an open problem.

The existence of a unique minimizer of the associated energy functional

$$I(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2} \right) dx - \sum_{i=1}^k a_i \phi(x_i),$$

follows from Theorem 1.3. As before, we denote this minimizer by ϕ_ρ . Our concern in this section consists in proving that this minimizer solves (4.1) in a weak or a strong sense. We can prove this fact in some particular cases only.

We need some intermediate steps. First we prove that ϕ_ρ satisfies strongly (4.1) in $\mathbb{R}^N \setminus \Gamma$, where

$$\Gamma = \bigcup_{i \neq j} \overline{x_i x_j}$$

and this is true without any restriction on the coefficients a_i and the location of the charges. This argument is already included in [18]. We give it for completeness.

Lemma 4.1. *The minimum ϕ_ρ of I satisfies strongly*

$$\begin{cases} -\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi|^2}} \right) = 0, & \text{in } \mathbb{R}^N \setminus \Gamma, \\ \phi(x) \rightarrow 0, & \text{as } x \rightarrow \infty. \end{cases}$$

Furthermore, we have that

- (i) $\phi_\rho \in C^\infty(\mathbb{R}^N \setminus \Gamma) \cap C(\mathbb{R}^N)$;
- (ii) ϕ_ρ is strictly spacelike on $\mathbb{R}^N \setminus \Gamma$;
- (iii) for $i \neq j$, either ϕ_ρ is a classical solution on $\overline{x_i x_j}$, or

$$\phi_\rho(tx_i + (1-t)x_j) = t\phi_\rho(x_i) + (1-t)\phi_\rho(x_j), \quad 0 < t < 1.$$

Proof. Let Ω be an arbitrary bounded open domain with smooth boundary in $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$. Here we repeat the same arguments of Subsection 3.2, using the same notations. The main difference now is that $\rho = 0$, since $\Omega \subset \mathbb{R}^N \setminus \{x_1, \dots, x_k\}$. As in Subsection 3.2, we infer that ϕ_ρ is strictly spacelike and $Q^-(\phi_\rho) = 0$ in $\Omega \setminus K$ where K is defined in (3.6). Furthermore, we have

$$\phi_\rho(tx + (1-t)y) = t\phi_\rho(x) + (1-t)\phi_\rho(y),$$

for every $0 < t < 1$, where $x, y \in \partial\Omega$ are such that $|\phi_\rho(x) - \phi_\rho(y)| = |x - y|$ and $\overline{xy} \subset \Omega$. Again, if $K = \emptyset$, then ϕ_ρ is strictly spacelike.

We now show that K contains at most Γ . Assume by contradiction there exist $x, y \in \partial\Omega$ such that $x \neq y$, $\overline{xy} \subset \Omega$ and $|\phi_\rho(x) - \phi_\rho(y)| = |x - y|$ and such that the straight line spanned by \overline{xy} intersects Γ at a finite number of points (possibly zero). Without loss of generality, we can assume that $\phi_\rho(x) > \phi_\rho(y)$. It is easy to see, again, that for all $t \in (0, 1)$

$$(4.2) \quad \phi_\rho(tx + (1-t)y) = \phi_\rho(y) + t|x - y|.$$

Observe also that, since the line spanned by x and y intersects Γ at a finite number of points only, we can arbitrarily stretch Ω in the direction \overline{xy} to build new open sets $\Omega' \subset \mathbb{R}^N \setminus \{x_1, \dots, x_k\}$ with smooth boundaries and such that $\Omega \subset \Omega'$. Observe that $\phi_\rho|_{\Omega'}$ is a minimizer for $I_{\Omega'}$ in $C_{\phi_\rho}^{0,1}(\Omega')$ and, by [3, Theorem 3.2], we have that (4.2) holds for all $t \in \mathbb{R}$ such that $tx + (1-t)y \in \Omega'$. Now, we reach a contradiction with the boundedness of ϕ_ρ by choosing Ω' long enough in the direction \overline{xy} .

Arguing in a similar way, we see that on each edge of Γ , either $Q^-(\phi_\rho) = 0$ or the full edge belongs to K , namely (iii) holds.

Assertion (i) follows from [3, Remark p. 147], assertion (ii) from [3, Corollary 4.2]. \square

The next Lemma shows somehow the continuity of the minimizer with respect to the coefficients a_i , $i = 1, \dots, k$. As a consequence, when the coefficients are small, the minimizer ϕ_ρ is smooth on Γ .

Lemma 4.2. *For any $\varepsilon > 0$ there exists $\sigma > 0$ such that, if $\max_{i=1, \dots, k} |a_i| < \sigma$, then $\|\phi_\rho\|_\infty < \varepsilon$.*

Proof. Assume by contradiction that there exists $c > 0$ such that, for all $n \in \mathbb{N}$, $n \geq 1$, there exist $a_1^n, \dots, a_k^n \in \mathbb{R}$ with $\max_{i=1, \dots, k} |a_i^n| < 1/n$ and such that $\|\phi_{\rho_n}\|_\infty > c$, where $\rho_n = \sum_{i=1}^k a_i^n \delta_{x_i}$ and ϕ_{ρ_n} is the minimum of the functional associated with ρ_n . Since $\rho_n \rightarrow 0$ in \mathcal{X}^* and $-\sum_{i=1}^k \delta_{x_i} \leq \rho_n \leq \sum_{i=1}^k \delta_{x_i}$, by Theorem 5.3, we infer that $\phi_{\rho_n} \rightarrow \phi_0 = 0$ uniformly in \mathbb{R}^N and we reach a contradiction. \square

Finally, we show that the minimizer is bounded uniformly with respect to the coefficients and that the location x_i , $i = 1, \dots, k$, do not influence this bound. Again, as a consequence, the minimizer ϕ_ρ is smooth on Γ as soon as the charges are far from each other.

Lemma 4.3. *There exists $C = C(a_1, \dots, a_k) > 0$ such that, for all $x_i \in \mathbb{R}^N$, $i = 1, \dots, k$, $\|\phi_\rho\|_\infty < C$.*

Proof. Since $I(\phi_\rho) \leq 0$ and using (2.1) and the continuous embedding of \mathcal{X} into $L^\infty(\mathbb{R}^N)$, we have

$$c\|\phi_\rho\|_\infty^2 \leq c\|\nabla\phi_\rho\|_2^2 \leq \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla\phi_\rho|^2}\right) dx \leq \langle \rho, \phi_\rho \rangle \leq C(a_1, \dots, a_k)\|\phi_\rho\|_\infty.$$

\square

With these lemmas, we can complete the proof of Theorem 1.6.

Proof of Theorem 1.6. We start with the key argument to prove that ϕ_ρ is a weak solution outside $\{x_1, \dots, x_k\}$.

Claim 1: for every bounded domain Ω such that $\bar{\Omega} \subset \mathbb{R}^N \setminus \{x_1, \dots, x_k\}$, there exists a unique distributional solution $\bar{\phi}_\rho$ of the problem

$$\begin{cases} -\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi_\rho|^2}} \right) = 0, & \text{in } \Omega, \\ \phi = \phi_\rho, & \text{on } \partial\Omega. \end{cases}$$

The proof of this claim follows from a general result of Trudinger on divergence elliptic operators with measurable coefficients, see [25, Theorem 3.2]. The assumption [25, (3.21)] is clearly satisfied in our setting since the right-hand side of the equation is zero in Ω .

Claim 2: we have $\bar{\phi}_\rho = \phi_\rho$. By the arbitrariness of Ω in $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$, we deduce that $\bar{\phi}_\rho$ is the unique distributional solution of

$$\begin{cases} -\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi_\rho|^2}} \right) = 0, & \text{in } \mathbb{R}^N \setminus \{x_1, \dots, x_k\}, \\ \phi(x) \rightarrow 0, & \text{as } x \rightarrow \infty. \end{cases}$$

Lemma 4.1 and the uniqueness of $\bar{\phi}_\rho$ then imply that $\bar{\phi}_\rho = \phi_\rho$ a.e. in \mathbb{R}^N .

Proof of (1). Using the same arguments as in the proof of Lemma 4.1, we aim to prove that $K = \emptyset$. Assume by contradiction that $K \neq \emptyset$. Then there exist $x, y \in \partial\Omega$ such that $x \neq y$, $\overline{xy} \subset \Omega$ and $|\phi_\rho(x) - \phi_\rho(y)| = |x - y|$. Without loss of generality, we can assume that $\phi_\rho(x) > \phi_\rho(y)$. It is easy to see that for all $t \in (0, 1)$

$$\phi_\rho(tx + (1 - t)y) = \phi_\rho(y) + t|x - y|.$$

Two possibilities occur: either \overline{xy} intersects Γ in a finite number of points (possibly zero), or \overline{xy} intersects Γ in an infinite number of points. In the first case, we conclude as before. In the second case, without loss of generality, we can assume that \overline{xy} coincide with a piece of $\overline{x_1x_2}$ and, applying if necessary again [3, Theorem 3.2], we can assume, moreover, that $\overline{xy} = \overline{x_1x_2}$. Fixing $\varepsilon > 0$ such that $2\varepsilon < \min_{i,j=1,\dots,k, i \neq j} |x_i - x_j|$, by Lemma 4.2, there exists $\sigma > 0$ such that, if $\max_{i=1,\dots,k} |a_i| < \sigma$, then $\|\phi_\rho\|_\infty < \varepsilon$. Since we can find $x', y' \in \overline{x_1x_2}$ with $|x' - y'| > 2\varepsilon$, we reach a contradiction, indeed

$$2\varepsilon < |x' - y'| = |\phi_\rho(x') - \phi_\rho(y')| < 2\varepsilon.$$

The behavior of the gradient of ϕ_ρ near the singularities x_i is a consequence of [11, Theorem 1.5] (see also [18, Theorem 1.4]).

Proof of (2). This simply follows by modifying the arguments used to prove (1) and taking $\tau = 2C$, where C is given by Lemma 4.3. \square

Remark 4.4. Since ρ is the divergence of the field

$$F(x) = \sum_{i=1}^k b_i \frac{x - x_i}{|x - x_i|^N},$$

where $b_i = \frac{a_i}{(N-2)|\mathbb{S}^{N-1}|}$, the equation (4.1) can be written

$$-\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi_\rho|^2}} \right) = \operatorname{div} F(x), \quad \text{for } x \in \mathbb{R}^N.$$

Trudinger's result [25, Theorem 3.2] then applies on the whole of \mathbb{R}^N if

$$\left| \sum_{i=1}^k \sum_{j=1}^k b_i b_j \int_{\Omega} \frac{(x - x_i) \cdot (x - x_j)}{|x - x_i|^N |x - x_j|^N} \sqrt{1 - |\nabla \phi_\rho|^2} dx \right| < +\infty.$$

But checking this last assumption is delicate as it requires a precise study of the behaviour of $|\nabla \phi_\rho|$ around the points x_i .

Remark 4.5. It is claimed in [18, Corollary 3.2] that ϕ_ρ is smooth on those edges of Γ whose endpoints x_i, x_j satisfy $a_i a_j > 0$. However, the maximum principle or comparison arguments, as proposed therein to prove that ϕ_ρ achieves extremal values at the points x_i, x_j require on the one hand that the equation is at least weakly satisfied on a bounded open set containing x_i and on the other hand, [17, Theorem 8.1] assumes that the elliptic operator has bounded coefficients. Moreover, we do not have sign informations for ϕ_ρ on the boundary of a neighborhood of x_i . As a first step towards such a reasoning, it would be interesting to study the qualitative behaviour of the solution of

$$-\Delta \phi - \frac{1}{2} \Delta_4 \phi = \sum_{i=1}^k a_i \delta_{x_i},$$

in \mathbb{R}^3 . As far as we know, there is no result in the literature containing the precise asymptotics around the points x_i showing that those are either local minima or local maxima according to the sign of the coefficient a_i . One can explicitly compute the solution with a single charge. But in presence of a superpositions of charges, the two different scales in the differential operator makes the analysis of the problem non standard. The analysis is even more complicated with the curvature operator on the left-hand side.

We conclude this section by showing that we can prove that ϕ_ρ satisfies classically (4.1) in $\mathbb{R}^N \setminus \{x_1, \dots, x_k\}$ under some symmetry assumptions. This is yet another case where indeed ϕ_ρ solves the partial differential equation outside the points $\{x_1, \dots, x_k\}$. Assume for simplicity that we have two charges with equal coefficients, namely

$$\rho = a(\delta_{x_1} + \delta_{x_2}).$$

By uniqueness of the minimizer and since the functional is now invariant under the orthogonal transformations that exchanges x_1 and x_2 , we infer that ϕ_ρ is symmetric and therefore cannot be affine with slope 1 on the segment $\overline{x_1 x_2}$. Therefore, the assertion (iii) of Lemma 4.1 allows to conclude. The same argument can be used when we have a symmetric configuration of charges with equal coefficient.

In some special situations, we can argue without assuming any symmetry to prove that the minimizer is not affine with slope 1 on some of the edges of Γ . As an example, take three charges located at x_1, x_2 and x_3 and suppose that those points are not colinear. One can order the value of ϕ_ρ and assume without loss of generality that $\phi_\rho(x_1) < \phi_\rho(x_2) < \phi_\rho(x_3)$. Then, ϕ_ρ cannot be affine with slope 1 on $\overline{x_1 x_2}$ and $\overline{x_2 x_3}$ since otherwise we have

$$\phi_\rho(x_3) - \phi_\rho(x_1) = \phi_\rho(x_3) - \phi_\rho(x_2) + \phi_\rho(x_2) - \phi_\rho(x_1) = |x_3 - x_2| + |x_2 - x_1| > |x_3 - x_1|.$$

Other similar situations can be ruled out with the same argument but this is clearly an incomplete and unsatisfactory approach towards the understanding of the general case.

5. APPROXIMATIONS OF THE MINIMIZER

In this section, we provide several ways to approximate the minimizer ϕ_ρ by a sequence of solutions of some approximating PDEs. It is unfortunately unclear that the sequence of PDEs

leads to (BT) at the limit except if ρ is smooth but this case does not require any approximation procedure since it is covered by Theorem 1.5.

5.1. Approximation through a finite order expansion of the Lagrangian. One way to overcome, in some sense, the non differentiability of the functional (1.13) was proposed by Fortunato, Orsina and Pisani [15] where the authors consider $N = 3$ and observe that for b large,

$$(5.1) \quad \mathcal{L}_{\text{BI}} = b^2 \left(1 - \sqrt{1 - \frac{|\nabla\phi|^2}{b^2}} \right) \sim \frac{|\nabla\phi|^2}{2} + \frac{|\nabla\phi|^4}{8b^2}.$$

Then, for every density $\rho \in L^1(\mathbb{R}^3)$, the Euler equation

$$-\operatorname{div} \left(\left(1 + \frac{1}{2b^2} |\nabla\phi|^2 \right) \nabla\phi \right) = \rho, \quad \text{in } \mathbb{R}^3,$$

has an unique finite energy solution [15]. This means somehow that if we substitute Maxwell's Lagrangian by the right-hand side of (5.1), the contradiction to the principle of finiteness of the energy disappears.

We extend the study of approximated solutions [15], see also [27], by looking for higher order expansions, assuming ρ is in the dual space of $D^{1,2}(\mathbb{R}^N) \cap D^{1,2n}(\mathbb{R}^N)$ for some $n \geq 1$. This include the case of a Radon measure. Setting $b = 1$, for $n \geq 1$, we define \mathcal{X}_{2n} as the completion of $C_c^\infty(\mathbb{R}^N)$ with respect to the norm defined by

$$\|\phi\|_{\mathcal{X}_{2n}}^2 := \int_{\mathbb{R}^N} |\nabla\phi|^2 + \left(\int_{\mathbb{R}^N} |\nabla\phi|^{2n} dx \right)^{1/n}.$$

Formally, the operator Q^- (defined in (1.10)) can be expanded as a sum of $2h$ -Laplacian, namely

$$Q^-(\phi) = - \sum_{h=1}^{\infty} \alpha_h \Delta_{2h} \phi,$$

where for all $h \geq 1$, $\alpha_h > 0$ (the exact values of the coefficient α_h are given in [18, 27], they are not important for our purpose) and $\Delta_{2h} \phi := \operatorname{div}(|\nabla\phi|^{2h-2} \nabla\phi)$. The curvature operator is formally the Gateaux derivative of the functional

$$\int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla\phi|^2} \right) dx = \int_{\mathbb{R}^N} \sum_{h=1}^{\infty} \frac{\alpha_h}{2h} |\nabla\phi|^{2h} dx,$$

where the power series in the right hand side converges pointwise when $|\nabla\phi(x)| \leq 1$. Assuming $\rho \in \mathcal{X}_{2n}^*$, let us denote the n th approximation of the functional (1.13) by

$$I_n := \phi \in \mathcal{X}_{2n} \mapsto \sum_{h=1}^n \frac{\alpha_h}{2h} \int_{\mathbb{R}^N} |\nabla\phi|^{2h} dx - \langle \rho, \phi \rangle_{\mathcal{X}_{2n}}.$$

This functional is C^1 and we have existence and uniqueness of a critical point.

Proposition 5.1. *Given $n_0 \geq 1$ and $\rho \in \mathcal{X}_{2n_0}^*$, then, for all $n \geq n_0$, the functional $I_n : \mathcal{X}_{2n} \rightarrow \mathbb{R}$ has one and only one critical point.*

Proof. The proof is standard and does not deserve many details. Existence follows from the direct method of the calculus of variation since this functional is bounded from below, coercive and weakly lower semicontinuous on \mathcal{X}_{2n} . Uniqueness follows from the strict convexity of I_n . \square

We now describe the densities covered by this statement. As soon as $2n > \max\{N, 2^*\}$, Sobolev inequality combined with Morrey inequality show that \mathcal{X}_{2n} is continuously imbedded in $C_0^{0,\beta_n}(\mathbb{R}^N)$,

with $\beta_n = 1 - \frac{N}{2n}$, where we recall that $u \in C_0^{0,\beta_n}(\mathbb{R}^N)$ if there exists $C > 0$ such that for every $x, y \in \mathbb{R}^N$,

$$|u(x) - u(y)| \leq C|x - y|^{\beta_n} \quad \text{and} \quad \lim_{|x| \rightarrow \infty} u(x) = 0.$$

It follows that if $2n > \max\{N, 2^*\}$, $\mathcal{X}_{2n} \subset C_0(\mathbb{R}^N)$ and therefore $(C_0(\mathbb{R}^N))^* \subset \mathcal{X}_{2n}^*$.

As first important examples, we cover the case $\rho \in L^1(\mathbb{R}^N)$ since the linear functional

$$\phi \mapsto \int_{\mathbb{R}^N} \rho \phi \, dx$$

is bounded on $C_0(\mathbb{R}^N)$ and the case $\rho = \sum_{i=1}^k a_i \delta_{x_i}$ since the linear functional

$$\phi \mapsto \sum_{i=1}^k a_i \phi(x_i)$$

is also bounded on $C_0(\mathbb{R}^N)$. In fact, we can cover the case of Radon measure. Indeed by Riesz-Markov-Kakutani Representation Theorem, see for instance [14], $(C_0(\mathbb{R}^N))^*$ can be identified with $\mathcal{M}(\mathbb{R}^N; \mathbb{R})$, the space of signed finite Radon measures (i.e. Borel regular measures which are finite on compact sets of \mathbb{R}^N). This means that if $\rho \in (C_0(\mathbb{R}^N))^*$, there exists a unique Radon measure μ such that

$$\langle \rho, \phi \rangle = \int_{\mathbb{R}^N} \phi \, d\mu, \quad \text{for all } \phi \in C_0(\mathbb{R}^N).$$

Observe also that we can also deal with $L_{\text{loc}}^1(\mathbb{R}^N)$ densities. Indeed, the dual spaces of $C_c(\mathbb{R}^N)$ and $C_0(\mathbb{R}^N)$ coincide in the sense that if $\rho \in (C_0(\mathbb{R}^N))^*$, then the restriction of ρ to $C_c(\mathbb{R}^N)$ is a linear bounded functional whereas if $\rho \in (C_c(\mathbb{R}^N))^*$, then it has a unique extension $\bar{\rho} \in (C_0(\mathbb{R}^N))^*$ such that $\langle \bar{\rho}, \phi \rangle = \langle \rho, \phi \rangle$ for $\phi \in C_c(\mathbb{R}^N)$ and the norms of ρ and $\bar{\rho}$ are equal.

By interpolation, if $\phi \in \mathcal{X}_{2n}$, then $\phi \in D^{1,q}(\mathbb{R}^N)$ for any $2 \leq q \leq 2n$. It follows that \mathcal{X}_{2n}^* also contains the weak divergence of any vector field $\xi \in L^{q'}(\mathbb{R}^N; \mathbb{R}^N)$, where $1/q + 1/q' = 1$ and $2 \leq q \leq 2n$, i.e. $\frac{2n}{2n-1} \leq q' \leq 2$. In particular, observe that if $\xi \in L^1(\mathbb{R}^N; \mathbb{R}^N)$, then $\text{div } \xi \in \mathcal{X}_{2n}^*$ but in general we cannot conclude that $\text{div } \xi \in \mathcal{X}_{2n}^*$ for a finite integer n . This implies that for $\xi \in L^1(\mathbb{R}^N; \mathbb{R}^N)$, we can take $\rho = \text{div } \xi$ in Theorem 1.3 but this case is not covered by Proposition 5.1 nor by Theorem 5.2.

Let $\rho \in \mathcal{X}_{2n_0}^*$ for some $n_0 \geq 1$ and, for all $n \geq n_0$, let ϕ_n be the unique solution of the approximated problem given by Proposition 5.1. The next theorem was basically the heart of the existence of a minimizer of the functional in [18]. We give it here in a general setting, providing a detailed proof for completeness. We mainly follow the idea of [18, Subsection 2.4 (first part)].

Theorem 5.2. *If $\rho \in \mathcal{X}_{2n_0}^*$ for some $n_0 \geq 1$, then ϕ_n tends to ϕ_ρ weakly in \mathcal{X}_{2m} for all $m \geq n_0$.*

Proof. Let $\mathcal{I}_n = I_n(\phi_n)$. For every $m < n$, we have $\mathcal{I}_m \leq I_m(\phi_n) < \mathcal{I}_n$. Thus $(\mathcal{I}_n)_n$ is a strictly increasing sequence. Since $\mathcal{I}_n < 0$ for all n , such a sequence is bounded from above. Hence we infer that $\mathcal{I} := \lim_{n \rightarrow \infty} \mathcal{I}_n \in \mathbb{R}$.

Since, for all $m < n$, $I_m(\phi_n) < \mathcal{I}_n < \mathcal{I}$, we deduce that $\|\nabla \phi_n\|_{2m} \leq C(m)$ by coercivity. Then, we infer that $(\phi_n)_n$ is weakly convergent in \mathcal{X}_{2m} for all $m \geq 1$. Hence, by a diagonal argument, we conclude that the limit $\bar{\phi}$ is the same for all m and belongs to the $\bigcap_{m \geq 1} \mathcal{X}_{2m}$.

For every $m, n \geq 1$, with $n_0 \leq m \leq n$, since $\mathcal{X}_{2n} \subset \mathcal{X}_{2n_0}$ and $\mathcal{I}_n < 0$, we have

$$\begin{aligned} \frac{\alpha_1}{2} \|\nabla \phi_n\|_2^2 + \frac{\alpha_{n_0}}{2n_0} \|\nabla \phi_n\|_{2n_0}^{2n_0} &\leq \sum_{h=1}^m \frac{\alpha_h}{2h} \|\nabla \phi_n\|_{2h}^{2h} \\ &\leq \sum_{h=1}^n \frac{\alpha_h}{2h} \|\nabla \phi_n\|_{2h}^{2h} \\ &\leq \|\rho\|_{\mathcal{X}_{2n_0}^*} \|\phi_n\|_{\mathcal{X}_{2n_0}}. \end{aligned}$$

It follows that $\|\phi_n\|_{\mathcal{X}_{2n_0}} \leq C(n_0)$ and

$$\sum_{h=1}^m \frac{\alpha_h}{2h} \|\nabla \phi_n\|_{2h}^{2h} \leq C(n_0).$$

We can now take the limit as $n \rightarrow +\infty$ in this inequality. Indeed, $\phi_n \rightharpoonup \bar{\phi}$ in $D^{1,2m}(\mathbb{R}^N)$ so that the weak lower semicontinuity of the norms gives

$$\sum_{h=1}^m \frac{\alpha_h}{2h} \|\nabla \bar{\phi}\|_{2h}^{2h} \leq C(n_0),$$

for every $m \geq n_0$. This in turn implies that

$$\sum_{h=1}^{+\infty} \frac{\alpha_h}{2h} \|\nabla \bar{\phi}\|_{2h}^{2h} \leq C$$

from which we deduce that

$$\limsup_{h \rightarrow \infty} \|\nabla \bar{\phi}\|_{2h} \leq 1.$$

We now claim that $|\nabla \bar{\phi}| \leq 1$ a.e. in \mathbb{R}^N . Indeed, assume by contradiction that there exist $\varepsilon > 0$ and $\Omega \subset \mathbb{R}^N$ with $|\Omega| \neq 0$ such that $|\nabla \bar{\phi}| \geq 1 + \varepsilon$ a.e. in Ω . Then, for every $h \geq 1$

$$|\Omega|^{1/2h} (1 + \varepsilon) \leq \left(\int_{\Omega} |\nabla \bar{\phi}|^{2h} dx \right)^{1/2h} \leq 1$$

which is a contradiction for h large enough.

From here, we argue as in [18, Subsection 2.5] to prove that $\bar{\phi} = \phi_\rho$. We already know that $\bar{\phi} \in \mathcal{X}$. We first show that

$$(5.2) \quad I(\bar{\phi}) = \mathcal{I}$$

and then that

$$(5.3) \quad \mathcal{I} = \min_{\mathcal{X}} I(\phi).$$

To get (5.2), observe that, since $\mathcal{X} \subset \mathcal{X}_{2m}$ for every $m \geq 1$, we have $I_m(\bar{\phi}) \in \mathbb{R}$ and

$$\lim_{m \rightarrow \infty} I_m(\bar{\phi}) = I(\bar{\phi}).$$

Moreover $\mathcal{I}_m \leq I_m(\bar{\phi}) < I(\bar{\phi})$ and so $\mathcal{I} \leq I(\bar{\phi})$. On the other hand, since I_m is weakly lower semicontinuous, we have

$$I_m(\bar{\phi}) \leq \liminf_{n \rightarrow \infty} I_m(\phi_n).$$

But, since $I_m(\phi_n) < \mathcal{I}$ for $m < n$, we have that $I_m(\bar{\phi}) \leq \mathcal{I}$ and so, passing to the limit, we get that $I(\bar{\phi}) \leq \mathcal{I}$.

Finally, to show (5.3), assume by contradiction that there exists $\tilde{\phi} \in \mathcal{X}$ such that $I(\tilde{\phi}) = \mathcal{I} - \varepsilon$

for $\varepsilon > 0$. Then $\mathcal{I}_m \leq I_m(\tilde{\phi}) < \mathcal{I} - \varepsilon$ which contradicts the fact that $\mathcal{I}_m > \mathcal{I} - \varepsilon$ for m large enough. \square

We end up this section by observing that other approximation schemes can be used. The truncation of the power series to a finite order gives a lower approximation of the action functional. Another truncation was successfully proposed in [4] to deal with a related problem and could have been used here as well. Let us set $a_0(s) = (1 - s)^{-1/2}$ for all $s < 1$. Then

$$I(\phi) = \frac{1}{2} \int_{\mathbb{R}^N} A_0(|\nabla \phi|^2) dx - \langle \rho, \phi \rangle,$$

where $A_0(t) = \int_0^t a_0(s) ds$. Take $\theta \in (0, 1)$ and define $a_\theta : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ by

$$a_\theta(s) = \begin{cases} a_0(s) & \text{for } s \in [0, 1 - \theta] \\ \gamma s^{n-1} + \delta & \text{for } s \in (1 - \theta, +\infty) \end{cases}$$

where γ and δ are chosen in such a way that a_θ is C^1 . The truncated functional $I_{\theta,n} : \mathcal{X}_{2n} \rightarrow \mathbb{R}$ defined by

$$I_{\theta,n}(\phi) := \frac{1}{2} \int_{\mathbb{R}^N} A_\theta(|\nabla \phi|^2) dx - \langle \rho, \phi \rangle,$$

where $A_\theta(t) = \int_0^t a_\theta(s) ds$ gives another lower estimate of $I(\phi)$. Then we can show that given $n \geq 1$ and $\rho \in \mathcal{X}_{2n}^*$, the functional $I_{\theta,n}$ has one and only one critical point which is a weak solution of

$$-\operatorname{div}(a_\theta(|\nabla \phi|^2) \nabla \phi) = \rho.$$

In this approach, n is fixed which makes the functional setting easier than in the finite order approximation of the power series. It is chosen in such a way that $\rho \in \mathcal{X}_{2n}^*$. Taking a sequence $\theta_k \rightarrow 1$, we can show that the sequence of minimizers of $I_{\theta_k,n}$ converge to the minimizer of I , giving yet another way to approach the minimizer by a sequence of solutions of approximating problems.

5.2. Approximation by mollification of the charge density. In the previous section, we have shown that we can approximate the minimizer by a sequence of solutions of a family of approximating problems. We next show that, under sufficient conditions, we can approximate the minimizer by a sequence of solutions of the Born-Infeld equation. The starting point is the fact that a smooth charge yields a smooth minimizer that solves the Euler-Lagrange equation associated to the minimization problem. The approximation is then obtained by a first mollification of the charge.

As before, we still denote by ϕ_ρ the minimum associated to $\rho \in \mathcal{X}^*$ and we recall Definition 2.11 that we will use.

Theorem 5.3. *Let $\rho \in \mathcal{X}^*$ and suppose that there exist $(\rho_n)_n \subset \mathcal{X}^*$ and $\tilde{\rho} \in \mathcal{X}^*$ such that $\tilde{\rho} \geq 0$, $\rho_n \rightarrow \rho$ in \mathcal{X}^* and $-\tilde{\rho} \leq \rho_n \leq \tilde{\rho}$. Then ϕ_{ρ_n} converges to ϕ_ρ weakly in \mathcal{X} and uniformly in \mathbb{R}^N .*

Proof. Observe that by Lemma 2.12, $\phi_{\tilde{\rho}} \geq 0$ and $\phi_{-\tilde{\rho}} \leq \phi_{\rho_n} \leq \phi_{\tilde{\rho}}$. It is easy to see that $\phi_{-\tilde{\rho}} = -\phi_{\tilde{\rho}}$ and so $|\phi_{\rho_n}| \leq \phi_{\tilde{\rho}}$. By Lemma 2.13, there exists $\bar{\phi} \in \mathcal{X}$ such that ϕ_{ρ_n} converges to $\bar{\phi}$ weakly in \mathcal{X} and uniformly on compact sets. This implies that $|\bar{\phi}| \leq 2\phi_{\tilde{\rho}}$. Thanks to this uniform decay at infinity, it is easy to see that $\phi_{\rho_n} \rightarrow \bar{\phi}$ uniformly on \mathbb{R}^N .

To show that $\bar{\phi} = \phi_\rho$, let us denote

$$I_\rho(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx - \langle \rho, \phi \rangle.$$

Since $\rho_n \rightarrow \rho$ in \mathcal{X}^* , $\phi_{\rho_n} \rightarrow \bar{\phi}$ on \mathcal{X} and the first term of I is convex, we infer that

$$I_\rho(\bar{\phi}) \leq \liminf_{n \rightarrow \infty} I_{\rho_n}(\phi_{\rho_n}).$$

Moreover, since $I_{\rho_n}(\phi_{\rho_n}) \leq I_{\rho_n}(\phi_\rho)$ for all $n \geq 1$, and

$$I_\rho(\phi_\rho) = \lim_{n \rightarrow \infty} I_{\rho_n}(\phi_\rho),$$

we have

$$\limsup_{n \rightarrow \infty} I_{\rho_n}(\phi_{\rho_n}) \leq \lim_{n \rightarrow \infty} I_{\rho_n}(\phi_\rho) = I_\rho(\phi_\rho)$$

and we conclude. \square

By Theorem 1.5 and Theorem 5.3, we get the following approximation of ϕ_ρ as uniform limit of smooth solutions of a sequence of approximated problems.

Corollary 5.4. *Let $\rho \in \mathcal{X}^*$ and suppose that there exist $(\rho_n)_n \subset \mathcal{X}^* \cap L_{\text{loc}}^\infty(\mathbb{R}^N)$ and $\tilde{\rho} \in \mathcal{X}^*$ such that $\rho_n \rightarrow \rho$ in \mathcal{X}^* and $-\tilde{\rho} \leq \rho_n \leq \tilde{\rho}$. Then the sequence $(\phi_{\rho_n})_n$ of (locally strictly) spacelike solutions of (BT) with ρ_n converges to ϕ_ρ weakly in \mathcal{X} and uniformly in \mathbb{R}^N .*

Remark 5.5. *If $(\rho_n)_n \subset L^p(\mathbb{R}^N)$, with $1 \leq p < +\infty$, is such that $\rho_n \rightarrow \rho$ in $L^p(\mathbb{R}^N)$, we can immediately conclude that ϕ_{ρ_n} converges to ϕ_ρ weakly in \mathcal{X} and uniformly in \mathbb{R}^N . In particular, for a datum $\rho \in L^p(\mathbb{R}^N)$, the approximating sequence $(\phi_{\rho_n})_n$, where $(\rho_n)_n$ is a standard sequence of mollifications of ρ , is made of smooth strictly spacelike solutions of (BT) with the data ρ_n .*

6. THE BORN-INFELD-KLEIN-GORDON EQUATION AND OTHER EXTENSIONS

Another interesting problem which involves the Born-Infeld theory appears when we couple a field, governed by the nonlinear Klein-Gordon equation, with the electromagnetic field whose Lagrangian density is given by (1.5) or (1.6), by means of the Weil covariant derivatives.

In the wake of [10, 21], Yu, in [26], deals with the system

$$(6.1) \quad \begin{cases} \operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi|^2}} \right) = u^2(\omega + \phi), & x \in \mathbb{R}^3, \\ \Delta u = (m^2 - (\omega + \phi)^2) u - |u|^{p-2} u, & x \in \mathbb{R}^3. \end{cases}$$

Fixing u in a convenient space of radial functions, Yu considers the functional

$$E_u(\phi) = \int_{\mathbb{R}^3} \left[\left(1 - \sqrt{1 - |\nabla \phi|^2} \right) + \omega u^2 \phi + \frac{1}{2} \phi^2 u^2 \right] dx,$$

and proves that E_u possesses a minimizer ϕ_u without proving that the minimum ϕ_u is a critical point of E_u . Then the second equation of (6.1) is solved with ϕ_u in place of ϕ . Yu's conclusion is then that (u, ϕ_u) is a solution of (6.1) in a generalized sense, meaning that the second equation is classically satisfied while ϕ_u is a minimizer of E_u .

Our aim here is to show that the minimizer of E_u , and of similar functionals, is actually a solution of the corresponding equation. This leads us to consider more general equations of the form

$$(6.2) \quad \begin{cases} -\operatorname{div} \left(\frac{\nabla \phi}{\sqrt{1 - |\nabla \phi|^2}} \right) + f(x, \phi) = 0, & x \in \mathbb{R}^N, \\ \phi(x) \rightarrow 0, & \text{as } |x| \rightarrow \infty. \end{cases}$$

We assume that $f : \mathbb{R}^N \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function such that

(f1) there exists $p \geq 2^* - 1$ such that for all $(x, t) \in \mathbb{R}^N \times \mathbb{R}$

$$|f(x, t)| \leq C|t|^p;$$

- (f2) $f(\cdot, t)$ is radially symmetric, for all $t \in \mathbb{R}$;
- (f3) the functional $I_F : \mathcal{X} \rightarrow \mathbb{R}$ defined by

$$I_F(\phi) = \int_{\mathbb{R}^N} \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx + \int_{\mathbb{R}^N} F(x, \phi) dx,$$

where $F(x, t) = \int_0^t f(x, s) ds$, has a nontrivial radial local minimum ϕ_f in \mathcal{X} .

Remark 6.1. *The existence of a local minimum of I_F in \mathcal{X} follows, for example, by standard assumptions such as the coercivity of I_F and the convexity of the function $F(x, \cdot)$.*

We can prove the following

Theorem 6.2. *Suppose that (f1)-(f3) hold, then ϕ_f is a nontrivial weak solution of (6.2).*

Proof. Let $\phi_f \in \mathcal{X}$ be a local minimum of I_F . Arguing as in Proposition 2.7, since the map $\phi \in \mathcal{X} \mapsto \int_{\mathbb{R}^N} F(x, \phi)$ is of class C^1 , we infer that

$$\int_{\mathbb{R}^N} \frac{|\nabla \phi_f|^2}{\sqrt{1 - |\nabla \phi_f|^2}} dx + \int_{\mathbb{R}^N} f(x, \phi_f) \phi_f dx \leq 0$$

and that the set of points where $|\nabla \phi| = 1$ has zero measure. Then one concludes as in the proof of Theorem 1.4. \square

Remark 6.3. *Theorem 6.2 allows to complete the arguments of [26] concerning the existence of a nontrivial solution of (6.1). In that precise case, one can even conclude that the solution is classical and even smooth.*

We finish this section by showing that, if I_F has a nontrivial *local minimizer* in the generalized sense of Morse, i.e. the function is minimal with respect to compactly supported variations, see for example [1], then it is a solution of the (6.2). Of course, any local or global minimizer is a local minimizer in the sense of Morse. We emphasize that we do not require any radial symmetry here.

Theorem 6.4. *Suppose that (f1) holds and that there exists $\phi_0 \in \mathcal{X}$ such that for any bounded open set $\Omega \subset \mathbb{R}^N$, and for any $\phi \in \mathcal{X}$ with $\phi = \phi_0$ in $\mathbb{R}^N \setminus \Omega$, $I_F(\phi_0) \leq I_F(\phi)$. Then ϕ_0 is a weak solution of (6.2).*

Proof. Set $\rho(x) := f(x, \phi_0(x))$. For any $\Omega \subset \mathbb{R}^N$ bounded, arguing as in the previous sections, we infer that the functional I_Ω defined by

$$I_\Omega(\phi) = \int_{\Omega} \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx + \int_{\Omega} \rho \phi dx$$

has a unique minimizer $\psi_\Omega \in \mathcal{X}_\Omega$, where \mathcal{X}_Ω is the set of functions $\phi \in \mathcal{X}$ with $\phi = \phi_0$ in $\mathbb{R}^N \setminus \Omega$.

For every $x_0 \in \mathbb{R}^N$ and $R > 0$, we simply denote the minimizer $\psi_{B_{2R}(x_0)}$ by $\psi_{R,0}$ where $B_{2R}(x_0)$ is the ball of radius $2R$ centered at x_0 . We claim that

$$(6.3) \quad -\operatorname{div} \left(\frac{\nabla \psi_{R,0}}{\sqrt{1 - |\nabla \psi_{R,0}|^2}} \right) + f(x, \phi_0) = 0$$

on the ball $B_R(x_0)$ when $R > 0$ is sufficiently large. Since ϕ_0 is bounded, we may assume that $|\phi_0(x) - \phi_0(y)| \leq R/2$ for every $x, y \in \mathbb{R}^N$ and therefore, as soon as $|x - y| \geq R$, we have $|\phi_0(x) - \phi_0(y)| \leq |x - y|/2$. Arguing as Subsection 3.2, and keeping the same notations, we know that $\psi_{R,0}$ is strictly spacelike and solve (6.3) in $B_{2R}(x_0) \setminus K$ where

$$K = \{\overline{xy} \subset B_{2R}(x_0) \mid x, y \in \partial B_{2R}(x_0), x \neq y, |\phi_0(x) - \phi_0(y)| = |x - y|\}.$$

Of course, if $K \cap B_R(x_0) = \emptyset$, then our claim is proved. Assume by contradiction that $K \cap B_R(x_0) \neq \emptyset$. Then there exist $x, y \in \partial B_{2R}(x_0)$ such that $x \neq y$, $\overline{xy} \cap B_R(x_0) \neq \emptyset$ and $|\phi_0(x) - \phi_0(y)| = |x - y|$. But it is easy to see that in such a case $|x - y| \geq R$ and so $|\phi_0(x) - \phi_0(y)| \leq |x - y|/2$ which is a contradiction.

We now conclude by showing that $\psi_\Omega = \phi_0$ whatever $\Omega \subset \mathbb{R}^N$ which implies $\psi_{R,0} = \phi_0$ for every $x_0 \in \mathbb{R}^N$ and every $R > 0$. This follows from a totally standard argument in convex analysis. To simplify the notations, we set

$$J_\Omega(\phi) = \int_\Omega \left(1 - \sqrt{1 - |\nabla \phi|^2}\right) dx.$$

Since ϕ_0 is a local minimizer in the sense of Morse for I_F , we deduce that for $0 < t < 1$, we have

$$J_\Omega((1-t)\psi_\Omega + t\phi_0) + \int_\Omega F(x, (1-t)\psi_\Omega + t\phi_0) dx \geq J_\Omega(\phi_0) + \int_\Omega F(x, \phi_0) dx.$$

Using the convexity of J_Ω , we deduce that

$$(1-t)J_\Omega(\psi_\Omega) + tJ_\Omega(\phi_0) + \int_\Omega F(x, (1-t)\psi_\Omega + t\phi_0) dx \geq J_\Omega(\phi_0) + \int_\Omega F(x, \phi_0) dx,$$

or equivalently

$$J_\Omega(\psi_\Omega) \geq J_\Omega(\phi_0) + \frac{1}{t-1} \left(\int_\Omega F(x, (1-t)\psi_\Omega + t\phi_0) dx - \int_\Omega F(x, \phi_0) dx \right).$$

This yields

$$J_\Omega(\psi_\Omega) \geq J_\Omega(\phi_0) + \int_\Omega f(x, \phi_0(x))(\phi_0 - \psi_\Omega)$$

and therefore

$$I_\Omega(\psi_\Omega) \geq I_\Omega(\phi_0).$$

By uniqueness of the minimizer of I_Ω , we conclude that $\psi_\Omega = \phi_0$ in Ω . \square

REFERENCES

- [1] G. Alberti, L. Ambrosio, X. Cabré, *On a Long-standing Conjecture of E. De Giorgi: Symmetry in 3D for General Nonlinearities and a Local Minimality Property*, Acta Appl. Math. **65** (2001), 9–33. [24](#)
- [2] A. Azzollini, *Ground state solution for a problem with mean curvature operator in Minkowski space*, J. Funct. Anal., **266** (2014), 2086–2095. [6](#)
- [3] R. Bartnik and L. Simon, *Spacelike hypersurfaces with prescribed boundary values and mean curvature*, Comm. Math. Phys. **87** (1982), 131–152. [3](#), [6](#), [10](#), [14](#), [15](#), [16](#), [17](#)
- [4] D. Bonheure, A. Derlet, C. De Coster, *Infinitely many radial solutions of a mean curvature equation in Lorentz-Minkowski space*, Rend. Istit. Mat. Univ. Trieste **44** (2012), 259–284. [6](#), [22](#)
- [5] M. Born, *Modified field equations with a finite radius of the electron*, Nature **132** (1933), 282. [2](#)
- [6] M. Born, *On the quantum theory of the electromagnetic field*, Proc. Roy. Soc. London Ser. A **143** (1934), 410–437. [2](#)
- [7] M. Born, L. Infeld, *Foundations of the new field theory*, Nature **132** (1933), 1004. [3](#)
- [8] M. Born, L. Infeld, *Foundations of the new field theory*, Proc. Roy. Soc. London Ser. A **144** (1934), 425–451. [2](#), [3](#)
- [9] S.-Y. Cheng and S.-T. Yau, *Maximal space-like hypersurfaces in the Lorentz-Minkowski spaces*, Ann. of Math. **104** (1976), 407–419. [3](#)
- [10] P. d’Avenia, L. Pisani, *Nonlinear Klein-Gordon equations coupled with Born-Infeld type equations*, Electron. J. Differential Equations 2002, No. 26, 13 pp. [6](#), [23](#)
- [11] K. Ecker, *Area maximizing hypersurfaces in Minkowski space having an isolated singularity*, Manuscr. Math. **56** (1986), 375–397. [17](#)
- [12] I. Ekeland, R. Témam, *Convex analysis and variational problems*, SIAM, Philadelphia (1999). [4](#)
- [13] R.P. Feynman, R.B. Leighton, M. Sands, *The Feynman lectures on physics*, vol. 2, Addison-Wesley, London (1964). [2](#)
- [14] I. Fonseca, G. Leoni, *Modern Methods in the Calculus of Variations: L^p Spaces*, Springer, New York (2007). [20](#)

- [15] D. Fortunato, L. Orsina, L. Pisani, *Born-Infeld type equations for electrostatic fields*, J. Math. Phys. **43** (2002), 5698–5706. [2](#), [5](#), [19](#)
- [16] G.W. Gibbons, *Born-Infeld particles and Dirichlet p-branes*, Nuclear Phys. B **514** (1998), 603–639. [3](#)
- [17] D. Gilbarg, N.S. Trudinger, *Elliptic partial differential equations of second order*, Springer-Verlag, Berlin (2001). [18](#)
- [18] M.K.-H. Kiessling, *On the quasi-linear elliptic PDE $-\nabla \cdot (\nabla u / \sqrt{1 - |\nabla u|^2}) = 4\pi \sum_k a_k \delta_{s_k}$ in physics and geometry*, Comm. Math. Phys. **314** (2012), 509–523. [1](#), [5](#), [15](#), [17](#), [18](#), [19](#), [20](#), [21](#)
- [19] E.H. Lieb, M. Loss, *Analysis*, American Mathematical Society, Providence, RI, 2001. [2](#), [5](#)
- [20] J. Mawhin, *Nonlinear boundary value problems involving the extrinsic mean curvature operator*, Math. Bohem. **139** (2014), 299–313. [6](#)
- [21] D. Mugnai, *Coupled Klein-Gordon and Born-Infeld type equations: looking for solitary waves*, Proc. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci, **460** (2004), 1519–1527. [6](#), [23](#)
- [22] M.H.L. Pryce, *On a Uniqueness Theorem*, Math. Proc. Cambridge Philos. Soc. **31** (1935), 625–628. [3](#)
- [23] E. Serra, P. Tilli, *Monotonicity constraints and supercritical Neumann problems*, Ann. Inst. H. Poincaré Anal. Non Linéaire **28** (2011), 63–74. [11](#)
- [24] A. Szulkin, *Minimax principles for lower semicontinuous functions and applications to nonlinear boundary value problems*, Ann. Inst. H. Poincaré Anal. Non Linéaire **3** (1986), 77–109. [4](#), [8](#)
- [25] N.S. Trudinger, *Linear elliptic operators with measurable coefficients*, Ann. Scuola Norm. Sup. Pisa (3) **27** (1973), 265–308. [17](#), [18](#)
- [26] Y. Yu, *Solitary waves for nonlinear Klein-Gordon equations coupled with Born-Infeld theory*, Ann. Inst. H. Poincaré Anal. Non Linéaire **27** (2010), 351–376. [6](#), [23](#), [24](#)
- [27] F. Wang, *Solitary waves for the coupled nonlinear Klein-Gordon and Born-Infeld type equations*, Electron. J. Differential Equations 2012, no. 82, 12 pp. [6](#), [19](#)

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